

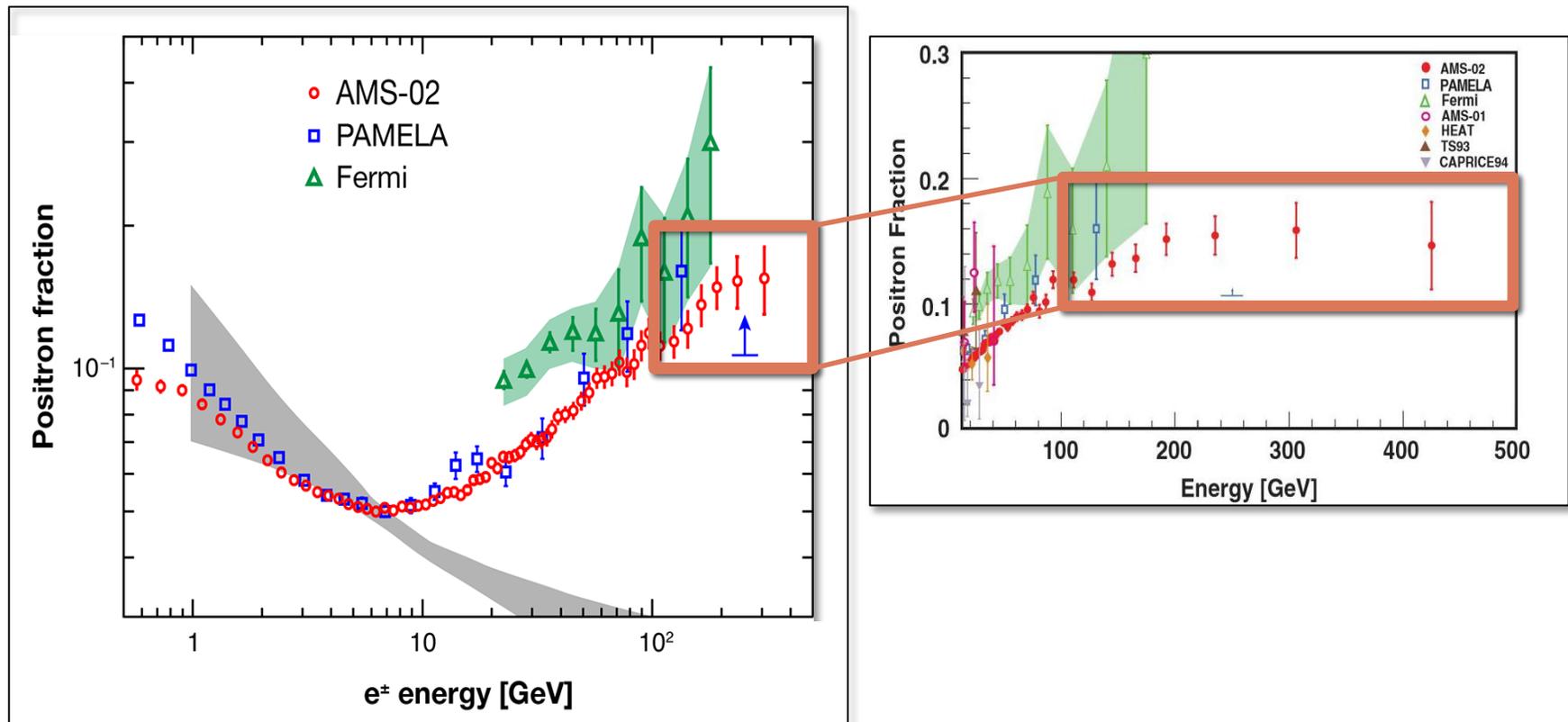


NEARBY PULSARS AND THE COSMIC RAY POSITRON EXCESS

Dan Hooper – Fermilab and the University of Chicago
IceCube Particle Astrophysics Symposium, Madison
May 9, 2017

The Cosmic Ray Positron Excess

- In 2008, PAMELA reported a surprisingly large quantity of positrons in the cosmic ray spectrum, now confirmed with much greater precision by AMS
- This result generated an explosive response from the particle dark matter community (~1800 citations, the majority of which focus on the implications for dark matter)

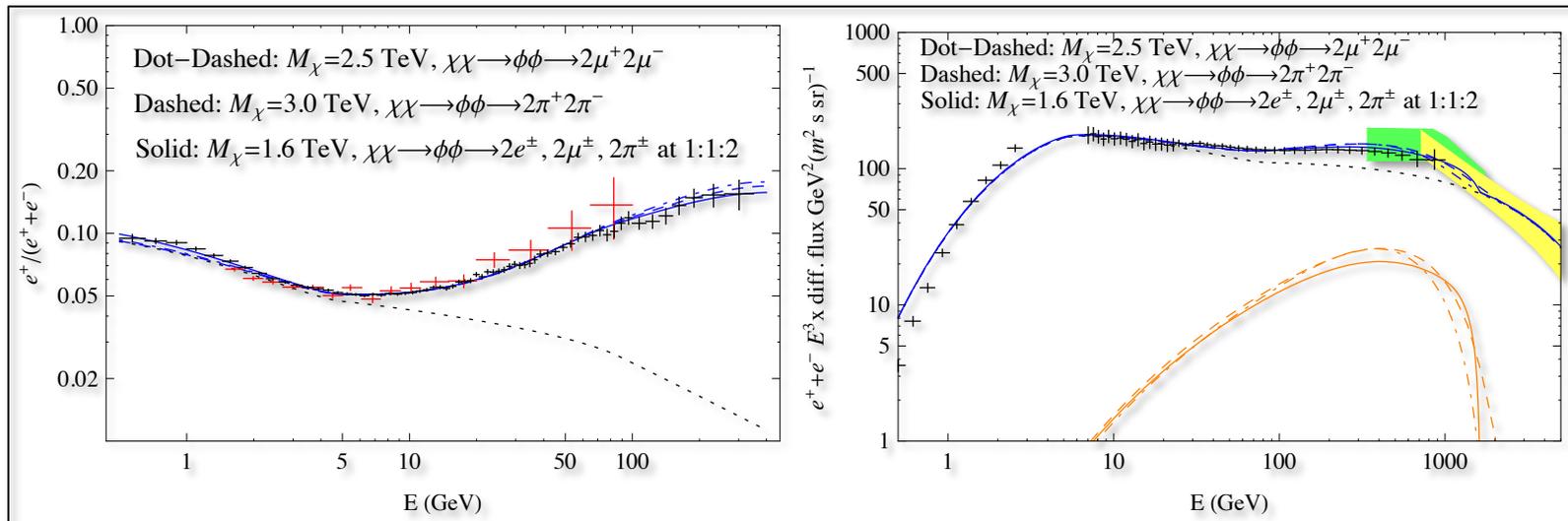


Where Do The Positrons Come From?

- The anticipated background to the positron flux is generated by cosmic ray interactions with gas in the ISM, yielding positrons through charged pion decay (*ie.* “*secondary*” positrons); this cannot account for the observed positrons
- Instead, three basic ideas have been proposed to account for the excess positrons:
 - 1) Annihilating or decaying dark matter particles
 - 2) The acceleration of secondary positrons within cosmic-ray sources (*ie.* supernova remnants)
 - 3) Nearby *primary* sources of high-energy positrons (*ie.* pulsars)

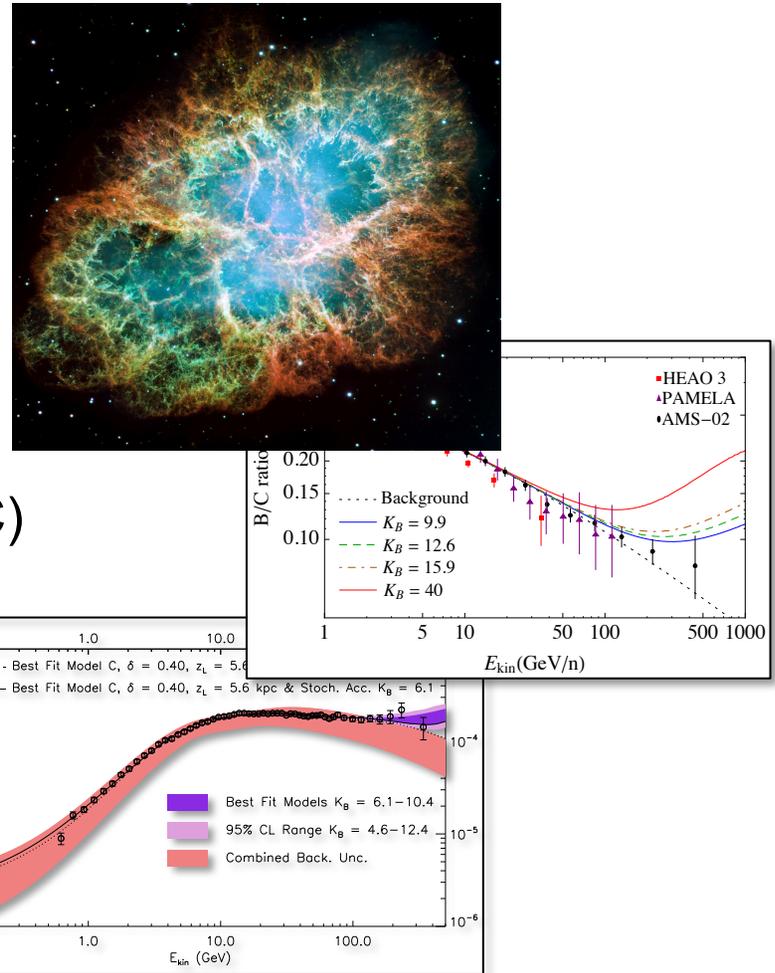
Annihilating Dark Matter and the Positron Excess

- In light of the detailed measurements of the positron fraction from AMS (and of the electron+positron spectrum from Fermi and HESS), few dark matter models can accommodate the data
- Dark matter models that *can* accommodate the data generally consist of a $\sim 1\text{-}3$ TeV particle that annihilates to unstable intermediate states, which then decay to electrons, muons and/or charged pions
- Large annihilation cross sections are also required ($\sim 10^{-24}$ to 3×10^{-23} cm³/s), making constraints from Fermi difficult to evade



The Acceleration of Secondary Positrons in Supernova Remnants

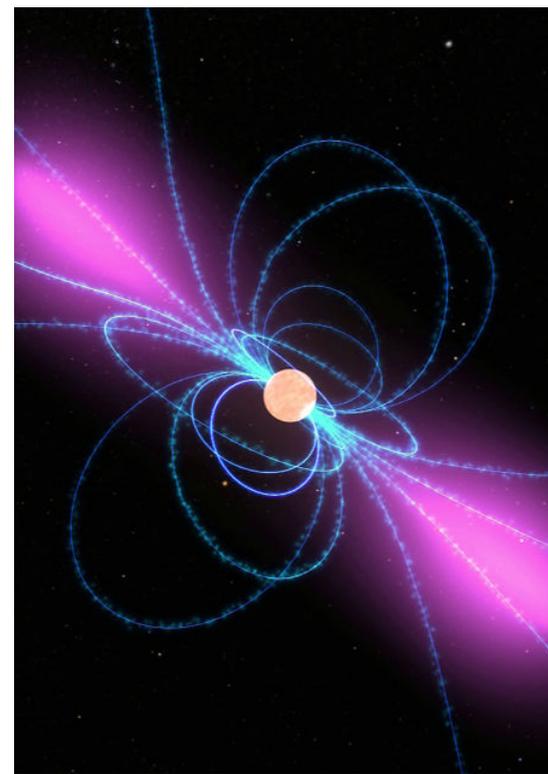
- Supernova remnants could generate secondary positrons and then accelerate them before they escape into the ISM
- If secondary positrons are accelerated in supernova remnants, then secondary antiprotons and boron nuclei should be accelerated as well
- Measurements of the boron-to-carbon (B/C) and antiproton-to-proton ratios from AMS indicate that secondary acceleration cannot account for the entirety of the positron excess, but may contribute non-negligibly



P. Blasi, PRL, arXiv:0903.2794; Mertsch, Sarkar, PRL, arXiv:0905.3152;
 Cholis, DH, PRD, arXiv:1312.2952; Cholis, DH, Linden, arXiv:1701.04406

Cosmic Ray Positrons From Pulsars

- Shortly after the PAMELA excess was reported, it was suggested that the positrons might originate from pulsars
- Pulsars are rapidly spinning neutron stars, which gradually convert their rotational kinetic energy into radio, X-ray, and gamma-ray emission, and into e^+e^- pairs
- Newly formed pulsars typically exhibit periods on the order of ~ 0.01 - 0.1 second, although most observed pulsars have higher periods (between ~ 0.1 and a few seconds)
- The rate of a pulsar's spin-down evolution, and its power depends on the strength of its magnetic field (which transfers rotational kinetic energy into radiation via magnetic dipole braking)



DH, Blasi, Serpico, PRD, arXiv:0810.1527;
Yuksel, Kistler, PRL, arXiv:0810.2784

(see also Zhang, Cheng, A&A, 2001; Grimani, A&A, 2007)

Pulsars Emission Models

- Considerable research activity has been focused on understanding exactly how pulsars generate their observed emission
- There are a number of basic elements that are found across a wide range of proposed models:
 - Electrons are accelerated by the strong magnetic fields, somewhere in the magnetosphere (the location is model dependent)
 - These electrons then induce electromagnetic cascades through the emission of curvature radiation
 - This results in the production of photons with energies above the threshold for pair production in the strong magnetic field
 - These electrons and positrons then escape the magnetosphere through open field lines, or after reaching the pulsar wind
- There is no consensus on what fraction of a pulsar's power is likely to go into the production of energetic e^+e^- pairs
- As high as ~20-30% of the energy budget? Or perhaps ~0.01%?

Which Pulsars Contribute to the Positron Flux?

Consider the standard cosmic-ray transport equation:

$$\frac{\partial}{\partial t} \frac{dn_e}{dE_e}(E_e, r, t) = \vec{\nabla} \cdot \left[D(E_e) \vec{\nabla} \frac{dn_e}{dE_e}(E_e, r, t) \right] + \frac{\partial}{\partial E_e} \left[\frac{dE_e}{dt}(r) \frac{dn_e}{dE_e}(E_e, r, t) \right] + \delta(r) Q(E_e, t)$$

Which Pulsars Contribute to the Positron Flux?

Consider the standard cosmic-ray transport equation:

$$\frac{\partial}{\partial t} \frac{dn_e}{dE_e}(E_e, r, t) = \vec{\nabla} \cdot \left[D(E_e) \vec{\nabla} \frac{dn_e}{dE_e}(E_e, r, t) \right] + \frac{\partial}{\partial E_e} \left[\frac{dE_e}{dt}(r) \frac{dn_e}{dE_e}(E_e, r, t) \right] + \delta(r) Q(E_e, t)$$

Diffusion: $D(E_e) = D_0 E_e^\delta$

Energy Losses: (ICS, Synchrotron)

$$-\frac{dE_e}{dt}(r) = \sum_i \frac{4}{3} \sigma_T \rho_i(r) S_i(E_e) \left(\frac{E_e}{m_e} \right)^2 + \frac{4}{3} \sigma_T \rho_{\text{mag}}(r) \left(\frac{E_e}{m_e} \right)^2$$

$$\equiv b(E_e, r) \left(\frac{E_e}{\text{GeV}} \right)^2$$

Injection Spectrum: (burst-like approximation)

$$Q(E_e, t) = \delta(t) Q_0 E^{-\alpha} \exp(-E_e/E_c)$$

Which Pulsars Contribute to the Positron Flux?

The solution to this equation is as follows:

$$\frac{dn_e}{dE_e}(E_e, r, t) = \frac{Q_0 E_0^{2-\alpha}}{8\pi^{3/2} E_e^2 L_{\text{dif}}^3(E_e, t)} \exp\left[\frac{-E_0}{E_c}\right] \exp\left[\frac{-r^2}{4L_{\text{dif}}^2(E_e, t)}\right]$$

where

$$L_{\text{dif}}(E_e, t) \equiv \left[\frac{D_0}{b(E_e/\text{GeV})^{1-\delta}(1-\delta)} \left(1 - (1 - E_e b t)^{1-\delta} \right) \right]^{1/2}$$

Taking the derivative of this solution with respect to r and setting it to zero, we find that a given pulsar will contribute the most to the local positron flux if it is located at a distance of $r \sim 2.4 L_{\text{dif}}$

Which Pulsars Contribute to the Positron Flux?

For parameters appropriate for the ISM: ($D_0 \simeq 2 \times 10^{28}$ cm²/s, $\delta \simeq 0.4$, $b = 1.8 \times 10^{-16}$ GeV/s).

$$L_{\text{dif}}(E_e, t) \simeq 200 \text{ pc} \left(\frac{35 \text{ TeV}}{E_e} \right)^{0.3} \left(1 - (1 - E_e b t)^{0.6} \right)^{1/2}$$

$$\sim 40 \text{ pc} \left(\frac{t}{10^5 \text{ yr}} \right) \left(\frac{E_e}{100 \text{ GeV}} \right)^{0.7}$$

Thus the pulsars that contribute the most to the local positron flux (those for which $r \sim 2.4 L_{\text{dif}}$) are those that are roughly $\sim 10^5$ years old *and* that located at a distance of roughly ~ 100 pc

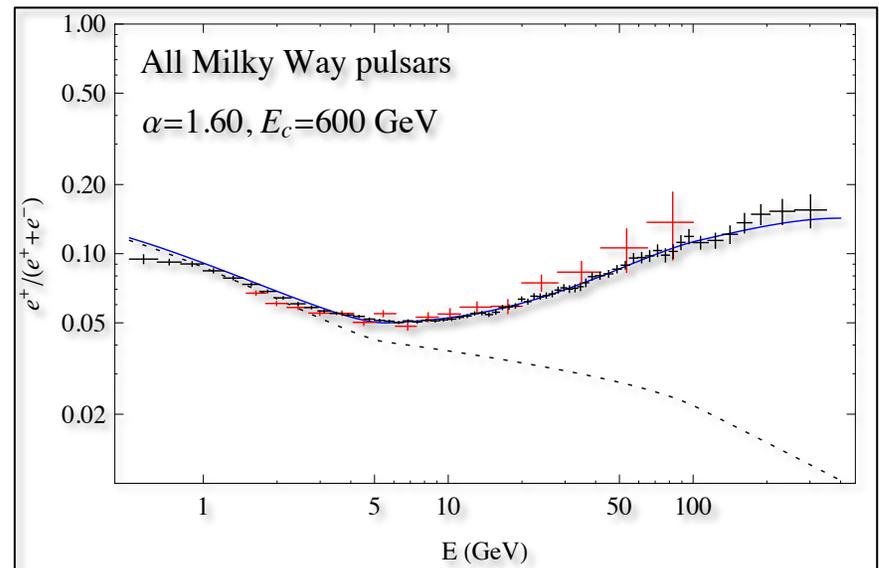
Cosmic Ray Positrons From Pulsars

- From these considerations, there are two known pulsars which stand out as the strongest potential sources of ~ 100 GeV cosmic ray positrons:

Geminga, age $\sim 370,000$ yrs, distance ~ 250 pc

B0656+14 (*ie.* monogem), age $\sim 110,000$ yrs, distance ~ 280 pc

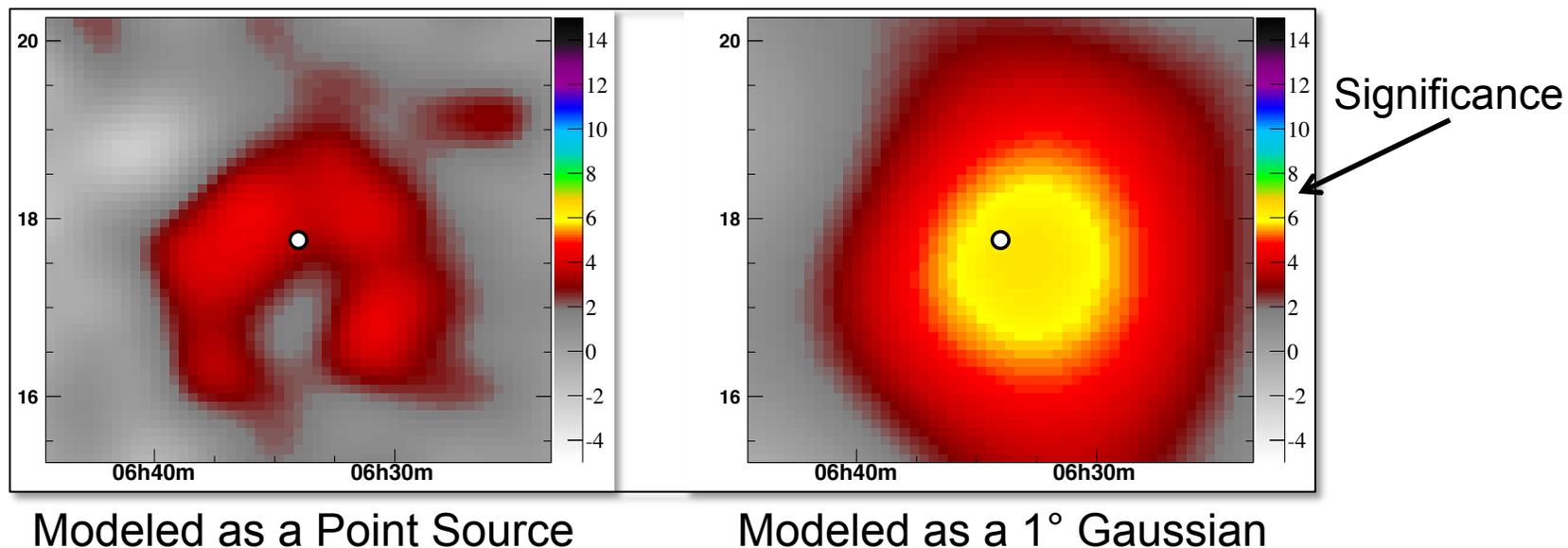
- If ~ 10 - 20% of the spin-down power of these pulsars is transferred into pairs, they could plausibly dominate the observed positron spectrum



DH, Blasi, Serpico, PRD, arXiv:0810.1527;
Yuksel, Kistler, PRL, arXiv:0810.2784;
Cholis, DH, PRD, arXiv:1304.1840

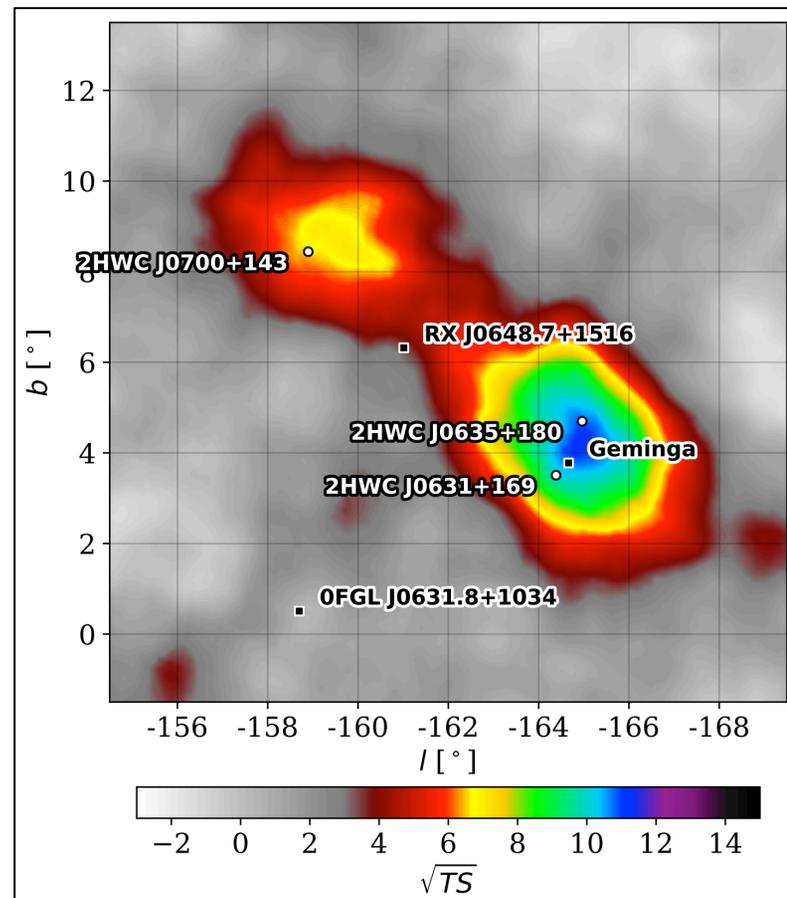
VHE Gamma-Ray Observations of Geminga

- Milagro has reported the detection of Geminga at an energy of ~ 35 TeV
- They also report the “definitive detection of extended emission” from Geminga, with a full-width-half-max of $2.6^{+0.7}_{-0.9}$ degrees



VHE Gamma-Ray Observations of Geminga

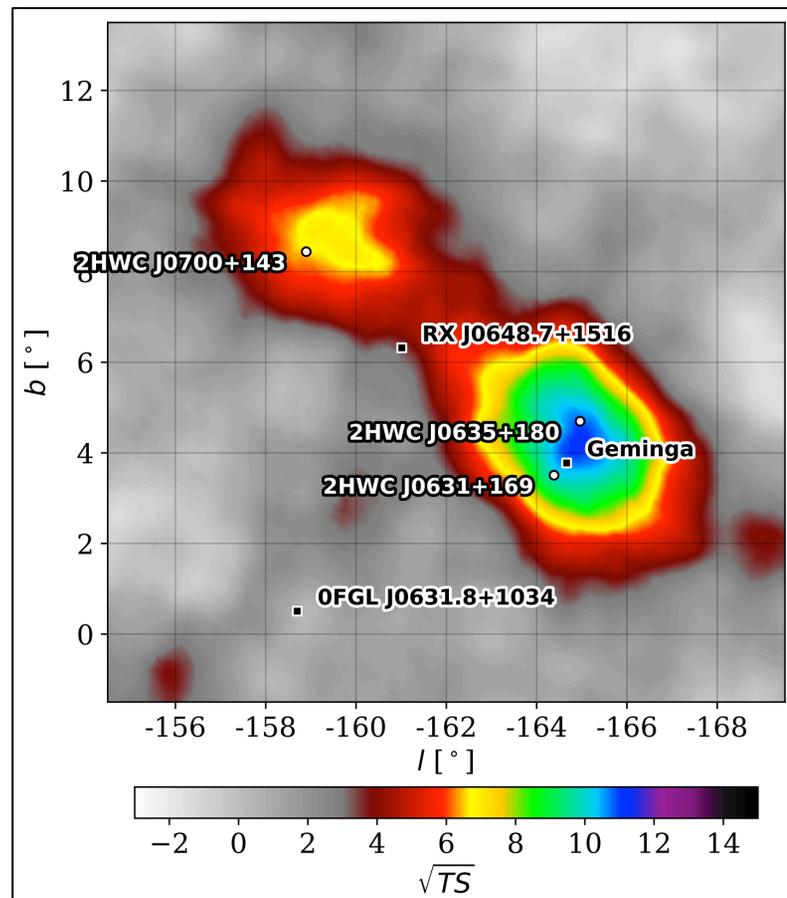
- Very recently, the HAWC Collaboration confirmed Milagro's detection of Geminga, and its spatial extension, in this case at ~ 7 TeV
- HAWC reports an extension of radius $\sim 2^\circ$, similar to that reported by Milagro
- Furthermore, HAWC also detects $\sim 2^\circ$ extended emission from the pulsar B0656+14 (2HWC J0700+143), not detected by Milagro (or by Fermi)



(Modeled as a 2° Radius Disk)

What Produces These Gamma Rays?

- The spatial extension of this emission indicates that the observed gamma rays do not originate from the pulsar itself, but from a region several parsecs in extent
- The only diffuse emission mechanisms that can produce such high-energy photons are inverse Compton scattering and pion production
- A pion production origin would require an implausibly large quantity of $\sim 10^2$ TeV protons ($>10^{46}$ erg), which would have to somehow be confined to the region for $>10^5$ years
- Inverse Compton scattering is almost certainly responsible for this emission



(Modeled as a 2° Radius Disk)

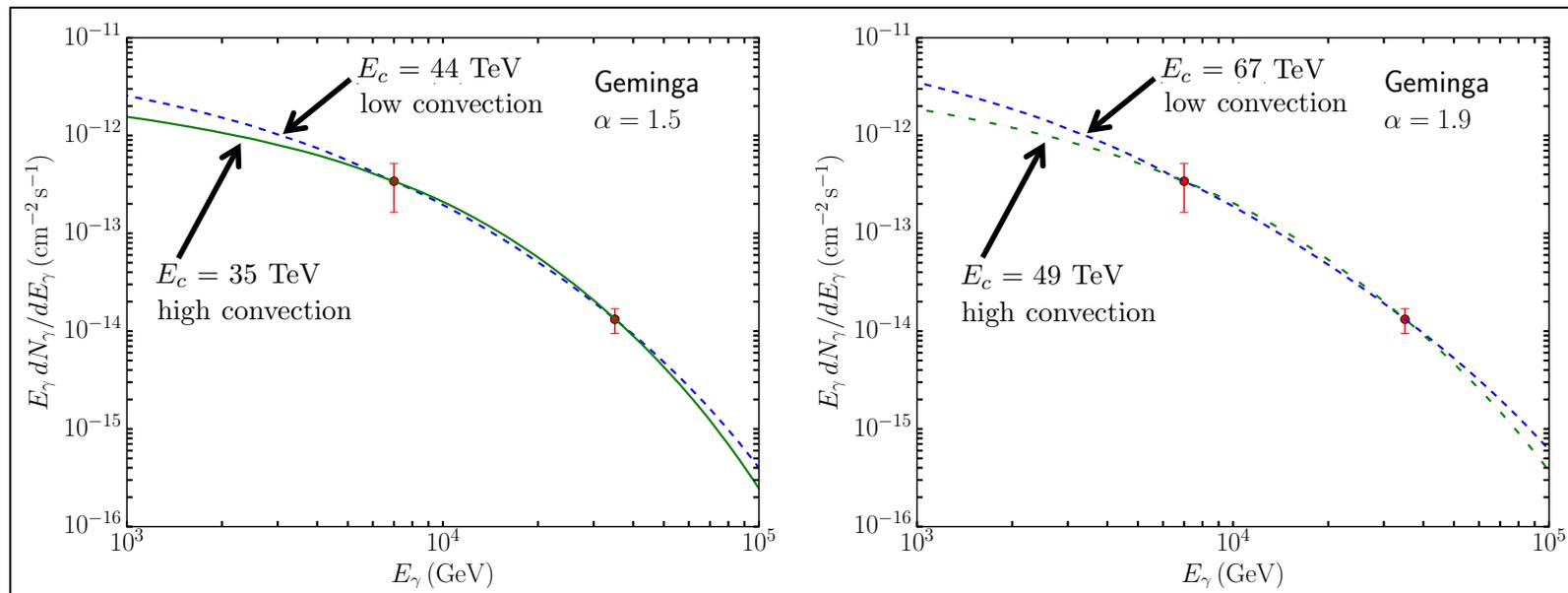
HAWC and Milagro Measurements Are Essential To Solving The Mystery Of The Positron Excess

- When a very high energy electron is injected into this environment, it emits the majority of its energy as Inverse Compton emission (along with a similar quantity as synchrotron)
- The results of HAWC and Milagro thus provide us with a direct measurement of the energy that Geminga and B0656+14 are currently injecting into very high-energy e^+e^- pairs (as well as information pertaining to the spectral shape of these pairs)

Main Idea: ***The spatial extension of Geminga and B0656+14 allow us to measure the critically important (and until now highly uncertain) fraction of these pulsars' spindown power that goes into the production of energetic e^+e^- pairs***

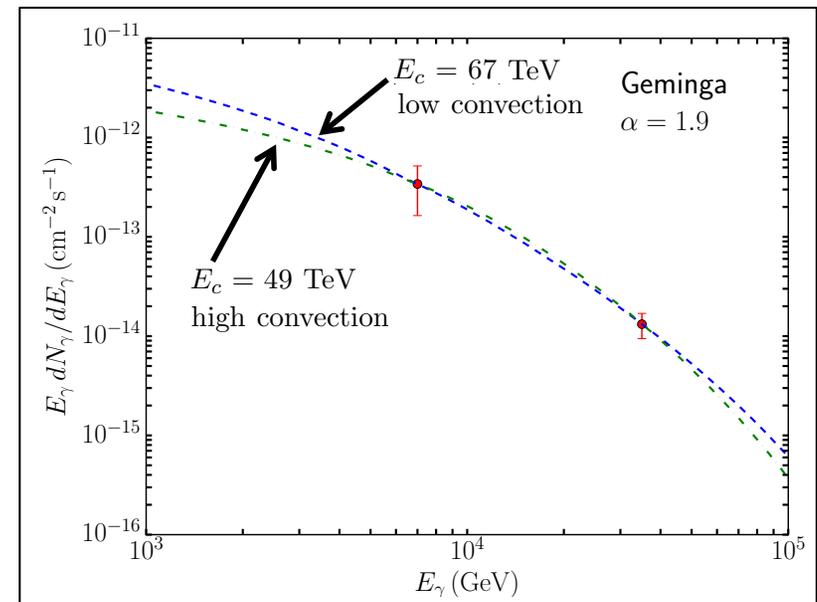
Implications of HAWC and Milagro for the Positron Excess

- For a given spectrum of injected pairs, we calculate the resulting ICS spectrum (including all Klein-Nishina corrections), and use this to constrain the normalization, spectral index (α), and energy cutoff (E_c)
- The VHE gamma-ray fluxes are best fit by $\alpha \sim 1.5-2.0$ and $E_c \sim 35-70$ TeV
- In these best-fit models, between 7-29% of Geminga's current spindown power goes into e^+e^- pairs – *similar to that required for the positron excess!*



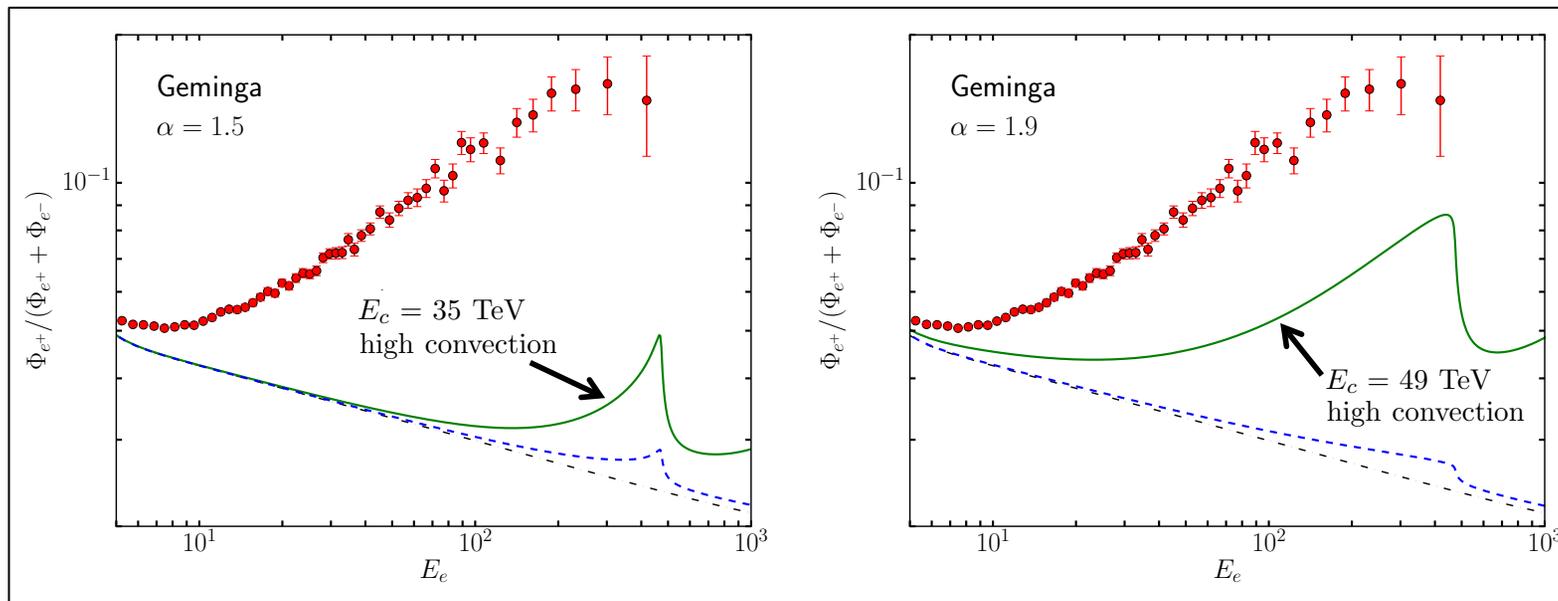
The Role Of Convection

- For the low degree of diffusion that is required to explain the observed extension, it is a combination of convection and energy-independent diffusion ($\delta \sim 0$) that enables lower energy electrons to escape the region surrounding Geminga – we parameterize the combination of these effects by a convection velocity
- The convection velocity impacts the shape of the gamma-ray spectrum, and when we take into account the slope reported by HAWC (-2.23 ± 0.08), we find that a sizable convection velocity is required $v_c \sim 100\text{-}500 \text{ km/s}$
- In these plots, “*high convection*” refers to $v_c \sim 230 \text{ km/s} \times (r_{\text{region}}/5 \text{ pc})$ – *focus on these curves*



Implications of HAWC and Milagro for the Positron Excess

- We can now use this information to calculate the contribution from Geminga to the local positron flux
- Across the range of models that provide a good fit to the HAWC and Milagro data, Geminga contributes non-negligibly to the observed excess



Some Caveats

Some Caveats

ICS vs synchrotron

- Some of the energy injected as e^+e^- pairs goes into synchrotron rather than ICS
- In our calculation, we adopted what we think are reasonable parameters ($B=3 \mu\text{G}$, $\rho_{\text{star}}=0.60 \text{ eV/cm}^3$, $\rho_{\text{IR}}=0.60 \text{ eV/cm}^3$, and $\rho_{\text{UV}}=0.10 \text{ eV/cm}^3$)
- If we had adopted a larger value of B , or smaller values of ρ_{star} , ρ_{IR} or ρ_{UV} , the contribution to the positron excess would increase (and vice versa)
- Over a reasonable range of these parameters, we could plausibly change the net result by up to a factor of roughly ~ 2 (either way)

Some Caveats

The time profile of Geminga's emission

- HAWC and Milagro measure the energy in ICS today, and thus are sensitive to the pairs that were injected in the past $\sim 10^4$ years
- In contrast, the positrons reaching the Solar System today were injected much longer ago, when the pulsar was young ($\sim 3 \times 10^5$ years ago)
- Geminga's rotation was faster and its spindown power higher when young:

$$\dot{E} = -\frac{8\pi^4 B^2 R^6}{3c^3 P(t)^4}$$

- In our calculation, we adopt the standard magnetic dipole braking model with $\tau \sim 10^4$ years:

$$P(t) = P_0 \left(1 + \frac{t}{\tau}\right)^{1/2} \quad \tau = \frac{3c^3 I P_0^2}{4\pi^2 B^2 R^6} \approx 9.1 \times 10^3 \text{ years} \left(\frac{P_0}{0.040 \text{ sec}}\right)^2$$

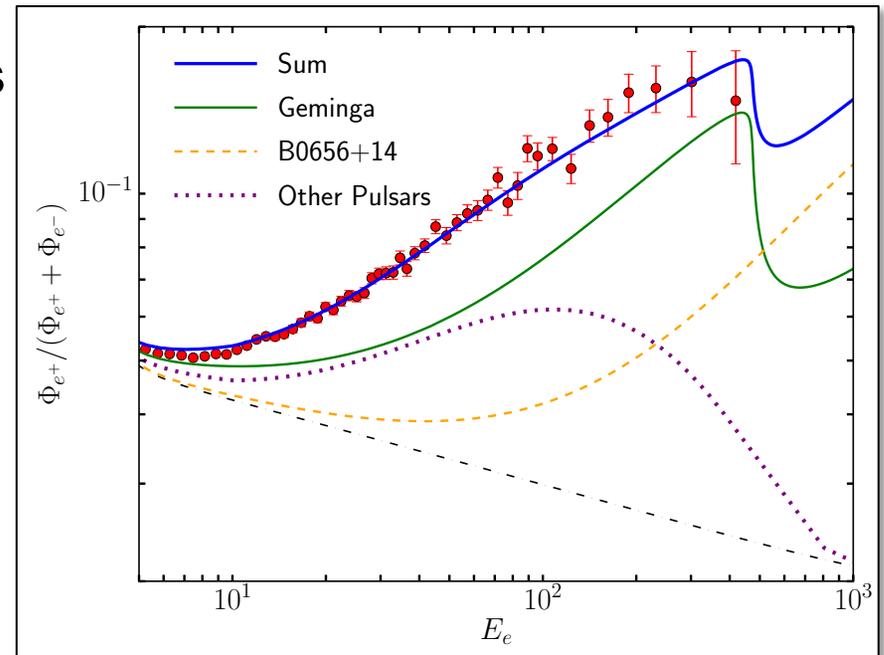
- By varying our choice of τ , we could plausibly change the net result by an order one factor

Positrons From Geminga, B0656+14, and More Distant Pulsars

- We have the most information about Geminga, and there is still an order one uncertainty as to its contribution to the local positron flux
- Larger uncertainties apply to B0656+14 and other pulsars
- That being said, can make a reasonable estimate for the total contribution

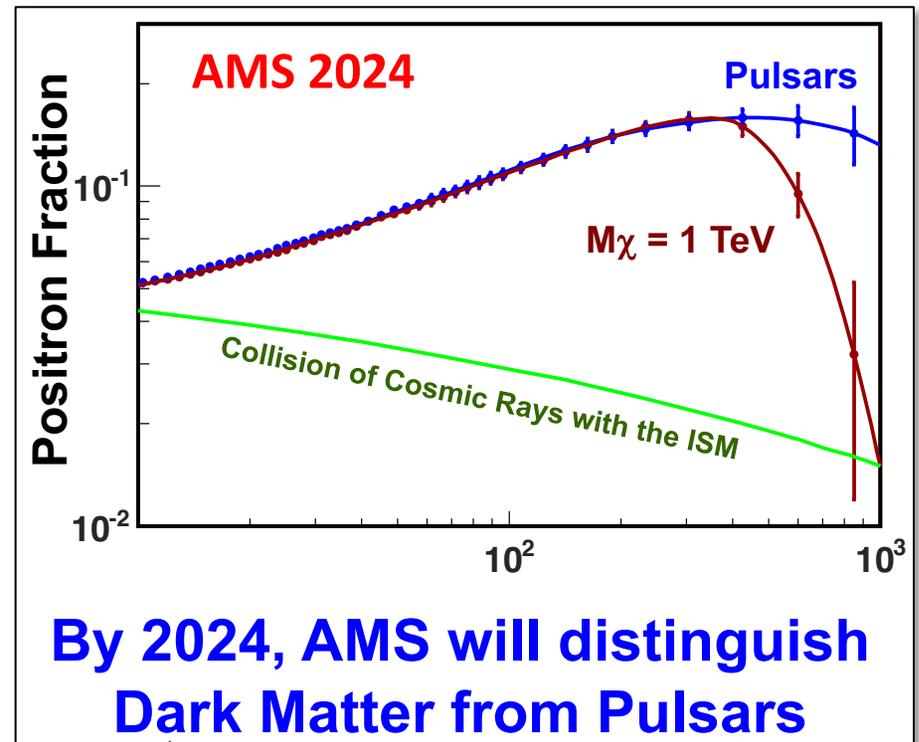
Positrons From Geminga, B0656+14, and More Distant Pulsars

- We have the most information about Geminga, and there is still an order one uncertainty as to its contribution to the local positron flux
- Larger uncertainties apply to B0656+14 and other pulsars
- That being said, can make a reasonable estimate for the total contribution
- In this figure, we have assumed that all pulsars inject e^+e^- pairs with the same efficiency and spectrum as Geminga, and adopted $\tau \sim 4.3 \times 10^3$ years and a birth rate of 2 new pulsars per century throughout the Milky Way (adopting the Lorimer *et al.* spatial distribution)
- These assumptions might not be precisely correct, but this shows that pulsars could very plausibly generate the entire excess, and likely provide the dominant contribution



A Note On Positron Spectral Features

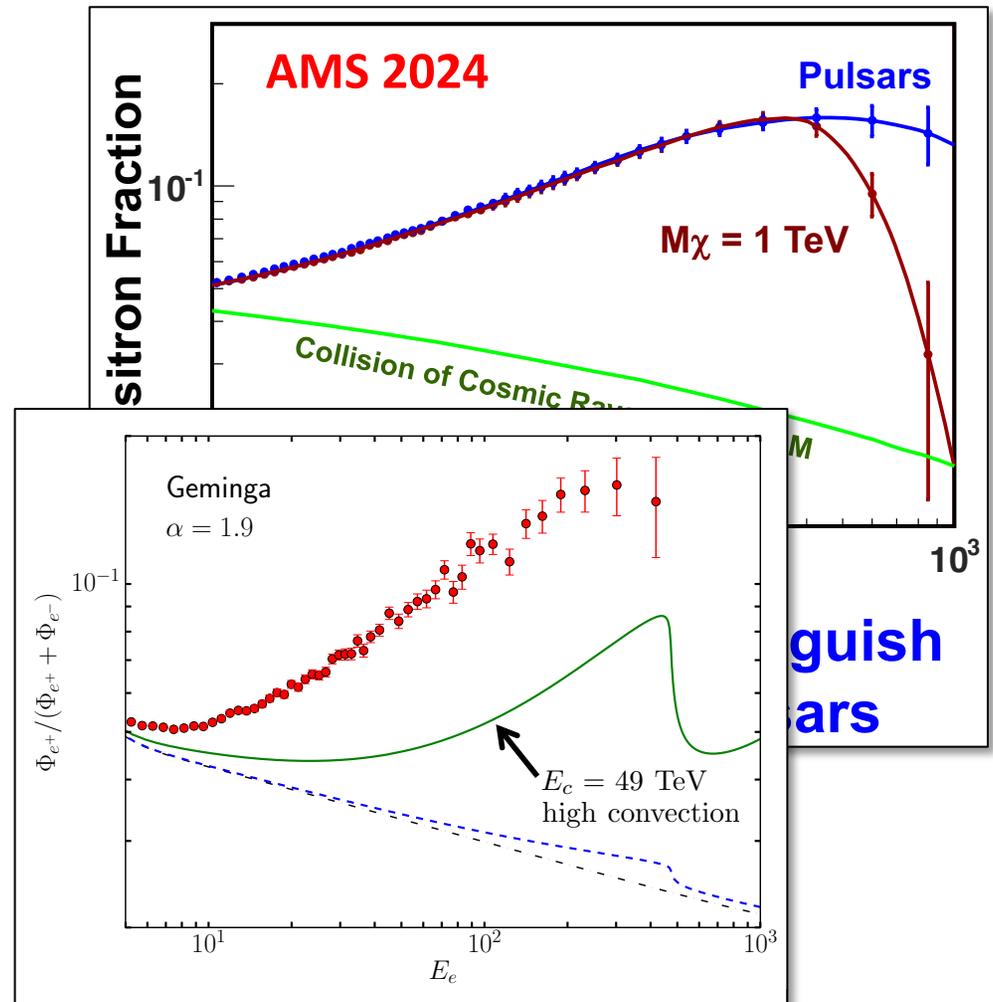
- A great deal is often made about “edges” and other spectral features that might appear in the positron spectrum
- Consider this plot, for example:



Sam's words, not mine

A Note On Positron Spectral Features

- A great deal is often made about “edges” and other spectral features that might appear in the positron spectrum
- Consider this plot, for example:
- A nearby pulsar could very plausibly generate an edge-like feature
- In fact, such an edge *will* appear at an energy of $E \sim I/bt_{age}$, which for Geminga is at $\sim 350\text{-}700$ GeV ($dE/dt = -bE^2$)



- Model based on J. Kopp PRD88, 2013

From Talk by Sam Ting, Dec. 2016

Summary

- Measurements from AMS-02 (as well as Fermi, HAWC) have revolutionized our understanding of cosmic rays in the Milky Way
- The PAMELA positron excess received a great deal of attention due to the possibility that it might be generated by annihilating dark matter – this no longer looks likely
- Recent observations of Geminga and B0656+14 by HAWC provide a determination of the flux of very high-energy e^+e^- pairs that is currently being injected by these sources – this efficiency factor was previously almost entirely unknown
- This new information implies that pulsars generate an order one fraction of the positron excess, and could very plausibly be responsible for the entirety of this signal

Personally, I think this is a very exciting result ... regardless of what Science Magazine has to say about it;)



Case weakens for antimatter sign of dark matter

By [Edwin Cartlidge](#) | Mar. 6, 2017 , 4:00 PM

A long debate over a mysterious surplus of antimatter—and whether it's a sign of dark matter—may be coming to **an anticlimactic end**. For more than a decade, multiple experiments have found an unexpected excess in the number of high-energy antielectrons, or positrons, in space, and some physicists suggested it could be due to particles of dark matter annihilating one another. Others countered with **a more mundane explanation**: The positrons come from rapidly rotating neutron stars, or pulsars. Now, a team of theorists has bolstered that more prosaic explanation, showing in detail that pulsars can indeed produce most or all of the excess.